

Measurement of time-dependent \mathcal{CP} violation in charmless B decays at LHCb

Stefano Perazzini on behalf of the LHCb Collaboration
 Istituto Nazionale di Fisica Nucleare (INFN)
 Sezione di Bologna
 Via Irnerio 46, 40126 Bologna, Italy

1 Abstract

In the following we present the measurements of time-dependent \mathcal{CP} violation in charmless B meson decays performed by LHCb analyzing the $p - p$ collision data collected at a center-of-mass energy of 7 TeV during the 2010 and 2011 LHC runs. In particular we will focus on the analysis of charmless two-body B decays where the direct and mixing-induced \mathcal{CP} asymmetry terms of the $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$ decays have been measured using 0.69 fb $^{-1}$ of data collected during 2011. The measurement of the branching ratio of the $B_s^0 \rightarrow K^{*0}\overline{K}^{*0}$ decay, using 35 pb $^{-1}$ collected during 2010, is also reported. In the end we show the relative branching ratios of all the decay modes of $B_{(s)}^0 \rightarrow K_S h^+ h'^-$ decays (where $h^{(\prime)} = \pi, K$), measured analyzing 1 fb $^{-1}$ of data collected during 2011.

2 Introduction

In this paper we report measurements of time-dependent \mathcal{CP} violation in charmless B meson decays performed at LHCb [1]. We will also briefly discuss the LHCb potential with other measurements in this sector. All the results here shown are obtained analyzing the $p - p$ collision data collected at a center-of-mass energy of 7 TeV during the 2010 and 2011 LHC runs. The time-dependent \mathcal{CP} asymmetry of either a B^0 or a B_s^0 meson decaying into a \mathcal{CP} -eigenstate f can be written as:

$$A_{\mathcal{CP}}(t) = \frac{\Gamma_{\overline{B} \rightarrow f}(t) - \Gamma_{B \rightarrow f}(t)}{\Gamma_{\overline{B} \rightarrow f}(t) + \Gamma_{B \rightarrow f}(t)} = \frac{A_{dir} \cos(\Delta m t) + A_{mix} \sin(\Delta m t)}{\cosh\left(\frac{\Delta \Gamma}{2} t\right) + A_{\Delta} \sinh\left(\frac{\Delta \Gamma}{2} t\right)}, \quad (1)$$

where $\Gamma(t)$ represents the time dependent decay rate of the initial B or \overline{B} meson to the final state f , Δm and $\Delta \Gamma$ are the B meson oscillation frequency and decay width difference respectively, and where the relation $A_{dir}^2 + A_{mix}^2 + A_{\Delta}^2 = 1$ holds. Within this parameterization, A_{dir} and A_{mix} account for \mathcal{CP} violation in the decay and in

the interference between mixing and decay, respectively. Since the decay amplitudes of charmless B decays receive important contribution from penguin diagrams, such a class of decays represents a powerful tool in the search for physics beyond the Standard Model. In fact new particles may appear as virtual contributions inside the loop diagrams, altering the Standard Model expectation for A_{dir} and A_{mix} .

The LHCb experiment has a great potential in this sector thanks to its capabilities of efficiently triggering and reconstructing charmless B decays, to its excellent decay-time resolution allowing to follow the fast B_s^0 oscillations and to the large cross section for B hadron production at LHC.

3 $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$

The $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$ decay amplitudes receive contributions from tree, electroweak penguin, strong penguin and annihilation topologies. The presence of penguin diagrams makes these decays sensitive to New Physics, but with the drawback of theoretical uncertainties in the determination of relevant hadronic parameters. In Refs. [3] strategies are presented in order to use the U -spin symmetry (*i.e.* the invariance of strong interaction dynamics under the quark exchange $d \leftrightarrow s$) relating the $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$ decays, in order to constraint hadronic parameters and determine the CKM phase γ and the B_s^0 mixing phase ϕ_s .

Crucial aspects for measuring time dependent \mathcal{CP} asymmetries in $B^0 \rightarrow \pi^+\pi^-$ and $B_s^0 \rightarrow K^+K^-$ decays are the event selection, the calibration of the particle identification (PID) and the flavour tagging. These decays are mainly selected by the hadronic trigger of LHCb, that exploits high transverse momentum and large impact parameter with respect to primary vertices, typical of B hadron decay products. Then the sample is further refined applying a set of kinematic and topological criteria. Finally the 8 final states ($\pi^+\pi^-$, K^+K^- , $K^+\pi^-$, π^+K^- , $p\pi^-$, $\bar{p}\pi^+$, pK^- and $\bar{p}K^+$) are separated into exclusive subsamples by means of the PID information provided by the two ring-imaging Cherenkov detectors of LHCb. The calibration of the PID observables is performed using large samples of pions, kaons and protons selected from $D^{*+} \rightarrow D^0(K^-\pi^+)\pi^+$ and $\Lambda \rightarrow p\pi^-$ decays (and their charge conjugates). As a demonstration of its PID capabilities, using 0.37 fb^{-1} of integrated luminosity collected during 2011, LHCb measured the relative branching ratios of various two-body charmless B decays [5]. The measurements of the branching ratios of the $B_s^0 \rightarrow K^+K^-$ and $B_s^0 \rightarrow \pi^+K^-$ decays are the most precise to date. In addition, the $B_s^0 \rightarrow \pi^+\pi^-$ decay is observed for the first time ever with a significance of more than 5σ .

The determination of the initial flavour of the signal B meson (the so-called "flavour tagging") is obtained using a multivariate algorithm that analyzes the decay products of the other B hadron in the event [6]. The response of the algorithm is

67 calibrated by measuring the oscillation of the flavour specific decay $B^0 \rightarrow K^+\pi^-$, in
 68 which the amplitude is related to the effective mistag rate. Using the measured mistag
 69 rate as an external input to a two dimensional (invariant mass and decay time) max-
 70 imum likelihood fit of the $\pi^+\pi^-$ and K^+K^- spectra, the direct and mixing-induced
 71 CP asymmetry terms $A_{\pi\pi}^{dir}$, $A_{\pi\pi}^{mix}$, A_{KK}^{dir} and A_{KK}^{mix} have been measured. The results,
 72 obtained using 0.69 fb^{-1} of data collected during 2011, are [7]:

$$A_{\pi\pi}^{dir} = 0.11 \pm 0.21(\text{stat.}) \pm 0.03(\text{syst.}) \quad (2)$$

$$A_{\pi\pi}^{mix} = -0.56 \pm 0.17(\text{stat.}) \pm 0.03(\text{syst.}) \quad (3)$$

$$A_{KK}^{dir} = 0.02 \pm 0.18(\text{stat.}) \pm 0.04(\text{syst.}) \quad (4)$$

$$A_{KK}^{mix} = 0.17 \pm 0.18(\text{stat.}) \pm 0.05(\text{syst.}), \quad (5)$$

73 with correlations $\rho(A_{\pi\pi}^{dir}, A_{\pi\pi}^{mix}) = -0.34$ and $\rho(A_{KK}^{dir}, A_{KK}^{mix}) = -0.10$. $A_{\pi\pi}^{dir}$ and $A_{\pi\pi}^{mix}$
 74 are measured for the first time at a hadronic machine and are compatible with the
 75 previous results from B -Factories [8]. The time-dependent \mathcal{CP} asymmetry terms A_{KK}^{dir}
 76 and A_{KK}^{mix} are measured for the first time ever and are compatible with theoretical
 77 expectations [4].

78 4 Other charmless decays

79 Using other charmless B decays where time-dependent \mathcal{CP} asymmetries can be mea-
 80 sured, LHCb already produced results. Using just 35 pb^{-1} collected during 2010,
 81 LHCb observed for the first time the $B_s^0 \rightarrow K^{*0}\bar{K}^{*0}$ decay [9]. This decay is gov-
 82 erned by pure penguin topologies offering an interesting potential in the quest for
 83 New Physics. It has been pointed out [10] that the time dependent \mathcal{CP} asymmetries
 84 in this decay can be used to measure the CKM phase γ and the B_s^0 mixing phase
 85 ϕ_s owing to the information provided by the U -spin related decay $B^0 \rightarrow K^{*0}\bar{K}^{*0}$.
 86 The signal yield obtained from the fit to the invariant mass distribution is equal to
 87 49.8 ± 7.5 candidates and the corresponding branching ratio is $\mathcal{BR}(B_s^0 \rightarrow K^{*0}\bar{K}^{*0}) =$
 88 $(2.81 \pm 0.46(\text{stat.}) \pm 0.45(\text{syst.}) \pm 0.34(f_s/f_d)) \times 10^{-5}$, where the last error comes
 89 from the knowledge of the b -quark hadronization probabilities.

90 Another important result is the measurement of the relative branching ratios
 91 of all the decay modes of $B_{(s)}^0 \rightarrow K_S h^+ h'^-$ decays (where $h^{(\prime)} = \pi, K$) [11]. The
 92 dominant topologies governing these decays are $b \rightarrow q\bar{q}s$ loop transitions (with $q =$
 93 u, d, s), where new particles in several extensions of the Standard Model may appear
 94 as virtual contributions. Several strategies have been proposed in order to extract
 95 information on the CKM phase γ and also on the B^0 and B_s^0 mixing phases ϕ_d and
 96 ϕ_s [12] from the time-dependent analysis of these decays in the Dalitz plane. As
 97 a first step, using the full 1 fb^{-1} collected during the 2011, LHCb measured the
 98 relative branching ratios of all the $B_{(s)}^0 \rightarrow K_S h^+ h'^-$ modes with respect to the well

Decay	$\mathcal{BR} \times 10^{-6}$
$\mathcal{BR}(B^0 \rightarrow K_S K^\pm \pi^\mp)$	$5.8 \pm 0.9 \pm 0.9 \pm 0.2$
$\mathcal{BR}(B^0 \rightarrow K_S K^+ K^-)$	$26.3 \pm 2.0 \pm 2.0 \pm 1.1$
$\mathcal{BR}(B_s^0 \rightarrow K_S \pi^+ \pi^-)$	$11.9 \pm 3.0 \pm 2.0 \pm 0.5$
$\mathcal{BR}(B_s^0 \rightarrow K_S K^\pm \pi^\mp)$	$97 \pm 7 \pm 10 \pm 4$
$\mathcal{BR}(B_s^0 \rightarrow K_S K^+ K^-)$	$4.2 \pm 1.5 \pm 0.9 \pm 0.2$

Table 1: Branching ratios of the $B_{(s)}^0 \rightarrow K_S h^+ h'^-$ decays measured by LHCb. The first error is statistical, the second is systematic and the third is due to the uncertainty on $\mathcal{BR}(B^0 \rightarrow K_S \pi^+ \pi^-)$ that is used as a normalization.

99 established $\mathcal{BR}(B^0 \rightarrow K_S \pi^+ \pi^-)$. The results are reported in Tab. 1 where the first
100 error is statistical, the second is systematic and the third is due to the uncertainty
101 on $\mathcal{BR}(B^0 \rightarrow K_S \pi^+ \pi^-)$. First evidence of the two decays $B_s^0 \rightarrow K_S \pi^+ \pi^-$ and $B_s^0 \rightarrow$
102 $K_S K^+ K^-$ is obtained with a significance of 4.3σ and 3.3σ respectively. Furthermore,
103 LHCb observed for the first time the $B_s^0 \rightarrow K_S K^\pm \pi^\mp$ decay.

104 References

- 105 [1] A. A. Alves *et al.*, *JINST* **3**, S08005 (2008).
- 106 [2] N. Cabibbo, *Phys. Rev. Lett.* **10**, 531 (1963); M. Kobayashi and T. Maskawa,
107 *Prog. Theor. Phys* **49**, 652 (1973).
- 108 [3] R. Fleischer, *Phys. Lett. B* **459**, 306 (1999); R. Fleischer, *Eur. Phys. J. C* **52**,
109 267 (2007); R. Fleischer and R. Knegjens, *Eur. Phys. J. C* **71**, 1532 (2011);
110 M. Ciuchini *et al.*, *JHEP* **1210**, 029 (2012).
- 111 [4] B. Adeva *et al.* [LHCb Collaboration], arXiv:0912.4179v3 [hep-ex]
- 112 [5] R. Aaij *et al.* [LHCb Collaboration], arXiv:1206.2794v1 [hep-ex].
- 113 [6] R. Aaij *et al.* [LHCb Collaboration], *Eur. Phys. J. C* **72**, (2012); R. Aaij *et al.*
114 [LHCb Collaboration], LHCb-CONF-2012-026.
- 115 [7] R. Aaij *et al.* [LHCb Collaboration], LHCb-CONF-2012-007.
- 116 [8] H. Ishino *et al.* [Belle Collaboration], *Phys. Rev. Lett.* **98**, 211801 (2007);
117 B. Aubert *et al.* [BaBar Collaboration], arXiv:0807.4226v2 [hep-ex].
- 118 [9] R. Aaij *et al.* [LHCb Collaboration], *Phys.Lett.* **B709**, 50 (2011).

- 119 [10] M. Ciuchini *et al.*, *Phys.Rev.Lett.* **100**, 031802 (2008); B. Bhattacharya *et al.*,
120 *Phys.Lett.* **B717**, 403 (2012); S. Descotes-Genon *et al.*, *Phys.Rev.* **D85**, 034010
121 (2012).
- 122 [11] R. Aaij *et al.* [LHCb Collaboration], LHCb-CONF-2012-023.
- 123 [12] M. Ciuchini *et al.*, *Phys.Rev.* **D74**, 051301 (2006); M. Gronau *et. al*, *Phys.Rev*
124 **D75**, 014002 (2007); M. Ciuchini *et al.*, *Phys.Lett.* **B645**, 201 (2007).